

8.47) a) Show that $\Theta = \sin \theta$ is a solution for the 2p states ($l=1, m=\pm 1$)

$$\frac{1}{\sin \theta} \frac{d}{d\theta} \left(\sin \theta \frac{d\Theta}{d\theta} \right) + \left(l(l+1) - \frac{m^2}{\sin^2 \theta} \right) \Theta \quad (8.65)$$

b) Show that the sum of the wave functions

$$\psi_{2,1,\pm 1}(r, \theta, \phi) = R_{2p}(r) \sin \theta e^{\pm i\phi}$$

is the 2p_x state, whereas the difference is 2ip_y.

a) For $l=1, m=\pm 1$, the DE becomes $[m^2 = (\pm 1)^2 = 1]$

$$\frac{1}{\sin \theta} \frac{d}{d\theta} \left(\sin \theta \frac{d\Theta}{d\theta} \right) + \left(2 - \frac{1}{\sin^2 \theta} \right) \Theta = 0$$

For $\Theta = \sin \theta, \Theta' = \cos \theta$

$$\frac{1}{\sin \theta} \frac{d}{d\theta} (\sin \theta \cos \theta) + \left(2 \sin \theta - \frac{1}{\sin \theta} \right) = 0$$

$$\frac{1}{\sin \theta} (\cos^2 \theta - \sin^2 \theta) + \left(2 \sin \theta - \frac{1}{\sin \theta} \right) = 0$$

Multiplying through by $\sin \theta$

$$\cos^2 \theta - \sin^2 \theta + 2 \sin^2 \theta - 1 = 0$$

$$\cos^2 \theta + \sin^2 \theta - 1 = 0$$

$$1 - 1 = 0 \quad \text{QED! } \sin \theta \text{ is a solution!}$$

b) Take the sum of $\psi_{2,1,\pm 1} = R_{2p}(r) \sin \theta e^{\pm i\phi}$

TABLE 8.2 gives $R_{2p}(r) = \frac{1}{\sqrt{24} a_B^3} r e^{-r/2a_B} = \frac{r}{\sqrt{24} a_B^3} e^{-r/2a_B}$

Thus

$$\psi_{2,1,\pm 1} = \frac{1}{\sqrt{24} a_B^3} r \sin \theta e^{-r/2a_B} e^{\pm i\phi}$$

Taking the sum

$$\psi_{2,1,1} + \psi_{2,1,-1} = \frac{1}{\sqrt{24} a_B^3} e^{-r/2a_B} r \sin \theta (e^{i\phi} + e^{-i\phi})$$



8.47) CONTINUED

NOTE THAT $\cos \phi = \frac{1}{2}(e^{i\phi} + e^{-i\phi})$

Thus

$$\psi_{2,1,1} + \psi_{2,1,-1} = \frac{1}{\sqrt{24a_B^5}} e^{-r/2a_B} \underbrace{2r \sin \theta \cos \phi}_{x = r \sin \theta \cos \phi!}$$

$$\left| \psi_{2,1,1} + \psi_{2,1,-1} = \frac{1}{\sqrt{24a_B^5}} e^{-r/2a_B} (2x) = 2\psi_{2x} \right| \text{ THE SUM GIVES } 2\psi_x$$

TAKING THE DIFFERENCE

$$\psi_{2,1,1} - \psi_{2,1,-1} = \frac{1}{\sqrt{24a_B^5}} e^{-r/2a_B} r \sin \theta (e^{i\phi} - e^{-i\phi})$$

NOTING THAT $\sin \phi = \frac{1}{2i}(e^{i\phi} - e^{-i\phi})$

$$\psi_{2,1,1} - \psi_{2,1,-1} = \frac{1}{\sqrt{24a_B^5}} e^{-r/2a_B} \underbrace{2i r \sin \theta \sin \phi}_y = r \sin \theta \sin \phi!$$

$$\left| \psi_{2,1,1} - \psi_{2,1,-1} = \frac{1}{\sqrt{24a_B^5}} e^{-r/2a_B} (2iy) = 2\psi_{2y} \right| \text{ THE DIFFERENCE GIVES } 2\psi_y$$

Thus, following TZB II's discussion on pages 276 & 277, the $2p$ wave functions are

$$\begin{aligned} \psi_{2p_z} &= A z e^{-r/2a_B} \\ \psi_{2p_y} &= A (2iy) e^{-r/2a_B} \\ \psi_{2p_x} &= A (2x) e^{-r/2a_B} \end{aligned}$$

Though $|\psi|^2$ in each Cartesian coordinate direction has two symmetric maxima on each axis (e.g. Fig 8.20), the overall probability distribution in all coordinate directions is spherically symmetric. This is expected for a purely radial potential.